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# Separation of colour degree of freedom from dynamics in a soliton cellular automaton 

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#### Abstract

We present an algorithm to reduce the coloured box-ball system, a onedimensional integrable cellular automaton described by motions of several colours (kinds) of balls, into a simpler monochrome system. This algorithm extracts the colour degree of freedom of the automaton as a word which turns out to be a conserved quantity of this dynamical system. It is based on the theory of crystal basis and in particular on the tensor products of $s l_{n}$ crystals of symmetric and anti-symmetric tensor representations.


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## 1. Introduction

The soliton cellular automaton by Takahashi and Satsuma [TS] is a discrete dynamical system related to the KdV equation. This automaton can be described by motions of a finite number of balls on an array of boxes, so it would have been called a box and ball system or a box-ball system for short. The original box-ball system in [TS] has only one kind of ball. Generalizations to the systems with several kinds (colours) of balls were introduced and studied [T, TNS, TTM]. We shall call these systems coloured systems, and the original one the (basic) monochrome system. Some years ago, a connection between these coloured systems and crystal basis theory [K1, K2] was found [FOY, HHIKTT]. By the algebraic (combinatorial) methods in the crystal basis theory, the scattering rules of solitons in the coloured systems [TNS] were explained.

In contrast to this success in the study on the scattering rules, there remain some problems on the construction of their conserved quantities, or of their general $N$-soliton solutions. There are studies on this subject for the coloured systems by analytical methods [TNS, TTM] or by combinatorial methods [FOY, F]. However the expressions for conserved quantities (or N -soliton solutions) in the former are rather complicated, and in both cases we did not know whether we had already obtained the list of all conserved quantities. It is clear that the difficulties of finding a simple and complete description of all conserved quantities are
due to the colour degree of freedom of these systems. In fact for the monochrome system the construction of all conserved quantities and the linearization of its dynamics (which are equivalent to the construction of general $N$-soliton solutions) have been done [TTS, Tkg] (but in the basic case, explained below).

In this paper we present an algorithm to separate this colour degree of freedom from dynamics in the coloured systems. The idea used here is an isomorphism between tensor products of crystals of symmetric and anti-symmetric tensor representations [HKOTY]. Our algorithm reduces a coloured system to a monochrome system, giving a word (finite sequence of colours of the balls) at the same time. We shall show that this word itself is a conserved quantity, i.e. it does not change under the time evolution of the automaton.

Throughout this paper we shall call the automata with all the cells (boxes) having capacity one basic systems [T, TNS, FOY], and those with cells of various capacities inhomogeneous systems [TTM, HHIKTT, F]. Our algorithm of separation of colour degree of freedom works in both cases. Clearly the former is a special case of the latter, so it is enough to give a proof for the latter case only. However we shall give a complete description of the proof for the basic case first and then generalize it to the inhomogeneous case, because we think it is more accessible for many readers.

The layout of this paper is as follows. In section 2 we introduce the basic coloured system and explain how the algorithm of separation of colour degree of freedom is conducted. We use a description by a carrier of balls which we call a decoding carrier. In section 3 we recall basic notions in the theory of $s l_{n}$ crystals and their tensor product decomposition by the LittlewoodRichardson (LR) rule. In section 4 we show explicit formulae for the isomorphism between tensor products of crystals, and explain the symmetric group generated by them. Here the decoding carrier is identified with an element of a crystal for the anti-symmetric representation. In section 5 we recall a description of coloured systems by means of the crystals, and give a proof for the separation of colour degree of freedom with the tools prepared in the previous sections. In section 6 we generalize the proof to that for the inhomogeneous system. In section 7 concluding remarks are given. Some examples and details of calculations are given in the appendices.

## 2. The coloured system and the decoding carrier

We recall the algorithm for the time evolution of the basic coloured system [T, TNS]. To adjust to the notations in crystal basis theory and Young tableaux, we call '(a box containing) a ball with index $i$ ' simply '(letter) $i(\geqslant 2$ )', and identify 'an empty box' with '(letter) 1 '. Fix an integer $n \geqslant 2$. At time $t$ we have an infinite sequence of letters $1,2, \ldots, n$. The numbers of $2, \ldots, n$ is finite. We denote this state by $\mathbf{p}$, and write the state at $t+1$ as $T(\mathbf{p})$, which means that we define a time-evolution operator $T$ that applies on $\mathbf{p}$.

Definition 2.1 ([T, TNS]). The time evolution operator $T$ is given by $T=K_{2} \ldots K_{n}$ where $K_{i}$ are the operators which work as

1. Move every letter i only once.
2. Exchange the leftmost $i$ with its nearest right 1 .
3. Exchange the leftmost $i$ among the rest of the $i$ 's with its nearest right 1 .
4. Repeat this procedure until all the i's are moved.

Remark 2.2. The numbering of balls is opposite to those in [T, TNS].

Example 2.3. Here is an example of the time evolution of the monochrome system.

$$
\begin{aligned}
& \mathrm{t}=0 \text {.... } 22222 \text {. . . . . } 222 \text {. } 2 \text {. . . . . . . . . . . } \\
& \mathrm{t}=1 \text {......... } 22222 \text {.... } 2.222 \\
& \mathrm{t}=2 \text {.............. } 2222 \text {. } 2 \text {... } 2222 \text {. . . . . } \\
& \text { t = } 3 \text {................... } 2 \text {. } 222 \text {. . . . } 22222 \text {. }
\end{aligned}
$$

We denoted letter 1 (empty box) by ' $\because$
Example 2.4. Here is an example of the time evolution of the coloured system which has four kinds of colours ' $2,3,4,5$ '.


We introduce another operator $T_{4}$. For this purpose we consider a carrier of balls which we call a decoding carrier. The capacity of this carrier is 2 , but it cannot be empty, so it always has at least one ball. It cannot have balls of the same kind at once. It loads and/or unloads balls when it passes by each box. The loading-unloading processes are depicted as follows:






For instance in the process $e$, a carrier containing (two balls with indices) $\alpha$ and $\beta$ comes from the left to a box containing $\gamma$. It picks up the $\gamma$, puts the $\alpha$ into the box and goes to the right. Note that we are assuming $\alpha, \beta, \gamma \geqslant 2$ here, and if we interpret a vacancy as letter ' 1 ', the processes $a$ and $b$ become special cases of $e$, and $c$ and $d$ become those of $f$ and $g$, respectively. We distinguished the processes $e$ and $g$ to adjust to the notations used later, although they are formally the same.

Now we define the operator $T_{\text {घ }}$. Let the decoding carrier always have a ' 2 ' at the beginning. The carrier runs from left to right along the automaton state $\mathbf{p}$, and changes it to another state which we call $T_{\mathrm{t}}(\mathbf{p})$. It puts the ' 2 ' into the automaton state and takes off a letter $\geqslant 2$ from the state.

Example 2.5. In the following picture the decoding carrier runs along the automaton state $55432 \ldots$, changes it into the state .55422. and takes off the letter 3.


Example 2.6. The following diagram shows how an automaton state (the first row in example 2.4) will be changed by applying $T_{\square}$ repeatedly:

| $s=0$ | 55432 . . . 542 . . . 2 |
| :---: | :---: |
| $s=1$ | 55422 . . . . 532 . . . 4 |
| $s=2$ | 55222 . . . . 432 . . 5 |
| $s=3$ | . 52222 . . . 543 . . . 2 |
| $\mathrm{s}=4$ | . . 22222 . . 554 . . . 3 |
| $s=5$ | . 22222 . . . . 552 |
| $s=6$ | . . 22222 . . . . 522 . 5 |
| $s=7$ | . 22222 . . . . . . 2225 |
| $s=8$ | . . . . 22222 . . . . . . 222 . |

Here $s$ is the number of times we have run the decoding carrier. The number attached at the end of each row (except the last one) shows the letter which will be taken off from that row. Note that the automaton state in the last row is equal to the first row in example 2.3.

Suppose, as in example 2.6, we could remove all the letters $\geqslant 3$ by applying $T_{\natural}$ finitely many times on $\mathbf{p}$. Let $\tilde{\mathbf{p}}$ be the state obtained from $\mathbf{p}$ after removing all the letters $\geqslant 3$. We denote by $\mathbf{y}$ the word made by these removed letters in reverse order $(\mathbf{y}=55432542$ for example 2.6 ). We may write symbolically

$$
\begin{equation*}
\mathbf{p}=\tilde{\mathbf{p}} \oplus \mathbf{y} \tag{3}
\end{equation*}
$$

In the following sections we shall show that for any state $\mathbf{p}$ we can decompose it into this form (3), and that under the time evolution $T$ the following relation holds:

$$
\begin{equation*}
T(\mathbf{p})=T(\tilde{\mathbf{p}}) \oplus \mathbf{y} \tag{4}
\end{equation*}
$$

We call this property of the coloured system a separation of colour degree of freedom from dynamics. For the above example we illustrate this property in appendix A.

Clearly this implies that the word $\mathbf{y}$ here is a conserved quantity of the coloured system. On the other hand the path $\tilde{\mathbf{p}}$ is regarded as a state for the monochrome system, and it has its own conserved quantities (and its dynamics can be linearized) [Tkg, TTS]. From (3) and (4) they are also conserved quantities of the original coloured system. In this way we can construct all conserved quantities (and can linearize the dynamics) of the (basic) coloured system.

## 3. Crystals and their tensor products

The dynamics of the coloured system in the previous section can be described by using $s l_{n}$ crystals. In this section we recall some properties of the $s l_{n}$ crystals in [KN]. The elements of the crystals $B_{\lambda}$ are given by semi-standard Young tableaux with any shape $\lambda$ and with letters $1, \ldots, n$. In particular we shall consider the crystals $B_{(\ell)}$ and $B_{(1,1)}$ which we call $B_{\ell}$ and $B_{\natural}$, respectively. They are given by

$$
\begin{align*}
& B_{\ell}=\left\{\left.\begin{array}{|l|l|l|l|}
\hline \alpha_{1} & \alpha_{2} & \cdots & \alpha_{\ell} \\
\hline
\end{array} \right\rvert\, 1 \leqslant \alpha_{1} \leqslant \alpha_{2} \leqslant \cdots \leqslant \alpha_{\ell} \leqslant n\right\}, \\
& B_{\natural}=\left\{\begin{array}{|l|l}
\hline \alpha & \mid 1 \leqslant \alpha<\beta \leqslant n \\
\hline \beta & 1 \leqslant \alpha,
\end{array}\right. \tag{5}
\end{align*}
$$

as sets. We shall mainly use these special types of crystals in the description of the coloured systems.

Remark 3.1. They are special cases of the crystals for rectangular shape Young tableaux. Such crystals can also be regarded as perfect crystals for affine Lie algebra $\widehat{s l}_{n}[\mathrm{KMN}, \mathrm{S}]$ and this fact may reduce some arguments below. In this paper we would rather avoid using this difficult concept of perfect crystals.

We recall basic notions in the theory of crystals. See [KN] for details. Let $I=\{1, \ldots, n-1\}$ be the index set and $B$ an $s l_{n}$ crystal. For any $i \in I$ there are maps $\tilde{f}_{i}$ from $B \sqcup\{0\}$ to $B \sqcup\{0\}$ and maps $\varepsilon_{i}$ and $\varphi_{i}$ from $B$ to $\mathbb{Z}_{\geqslant 0}$. (We shall omit descriptions of the maps $\tilde{e}_{i}$ which are basically defined as $\left(\tilde{f}_{i}\right)^{-1}$.) It is always assumed that $\tilde{f}_{i} 0=0$. For $B=B_{\ell}$ or $B_{\natural}$ their definitions are given as follows.

1. ( $B=B_{\ell}$ ) If $b \in B$ has at least one $i$ then $\tilde{f}_{i}(b)$ is that obtained by replacing the rightmost $i$ with $i+1$. Otherwise, $\tilde{f}_{i}(b)=0$. The maps $\varphi_{i}$ and $\varepsilon_{i}$ are given by $\varphi_{i}(b)=\#(i$ 's in $b)$ and $\varepsilon_{i}(b)=\#(i+1$ 's in $b)$.
2. $\left(B=B_{\sharp}\right)$ If $i$ appears in $b \in B$ and $i+1$ does not, then $\tilde{f}_{i}(b)$ is that obtained by replacing $i$ with $i+1$. Otherwise, $\tilde{f}_{i}(b)=0$. The maps $\varphi_{i}$ and $\varepsilon_{i}$ are given by

$$
\begin{aligned}
& \varphi_{i}(b)=[i \text { appears in } b \text { and } i+1 \text { does not }], \\
& \varepsilon_{i}(b)=[i+1 \text { appears in } b \text { and } i \text { does not }],
\end{aligned}
$$

where $[$ true $]=1,[$ false $]=0$.
Also the functions $\tilde{f}_{i}, \varphi_{i}, \varepsilon_{i}$ for more general $B_{\lambda}$ can be defined as in $[\mathrm{KN}]$ but we do not use their explicit forms. We simply note that there is a unique special element called the highest weight element in any $B_{\lambda}$, and every element of $B_{\lambda}$ can be obtained by applying some sequence of $\tilde{f}_{i}$ 's on it. The highest weight element in $B_{\lambda}$ is given by the Young tableau of shape $\lambda$ with all the letters in its first row are 1's, those in its second row are 2's and so on.

An important property of crystals is their tensor product structure. Let $\lambda_{1}$ and $\lambda_{2}$ be any two Young diagrams, and $B_{\lambda_{1}}$ and $B_{\lambda_{2}}$ the associated $s l_{n}$ crystals. As a set the tensor product $B_{\lambda_{1}} \otimes B_{\lambda_{2}}$ is simply a direct product. The actions of the maps $\tilde{f}_{i}$ on $B_{\lambda_{1}} \otimes B_{\lambda_{2}}$ are given by

$$
\tilde{f}_{i}\left(b \otimes b^{\prime}\right)=\left\{\begin{array}{lll}
\tilde{f}_{i} b \otimes b^{\prime} & \text { if } & \varphi_{i}(b)>\varepsilon_{i}\left(b^{\prime}\right)  \tag{6}\\
b \otimes \tilde{f}_{i} b^{\prime} & \text { if } & \varphi_{i}(b) \leqslant \varepsilon_{i}\left(b^{\prime}\right)
\end{array}\right.
$$

Here $0 \otimes b^{\prime}$ and $b \otimes 0$ should be understood as 0 . The tensor product has a finite decomposition

$$
\begin{equation*}
B_{\lambda_{1}} \otimes B_{\lambda_{2}} \simeq \bigoplus_{v} B_{v}^{\oplus m_{v}} \tag{7}
\end{equation*}
$$

which is described by the Littlewood-Richardson (LR) rule [N]. Here vs are distinct Young diagrams and $m_{\nu}(\geqslant 1)$ is the multiplicity of $B_{v}$. The $\simeq$ denotes the isomorphism of $s l_{n}$ crystals, i.e. it commutes with the actions of $\tilde{f}_{i}$. Tensor products of three or more crystals are defined again by using (6). To define them, such formulae as $\varepsilon_{i}\left(b \otimes b^{\prime}\right)=$ $\max \left(\varepsilon_{i}(b), \varepsilon_{i}(b)+\varepsilon_{i}\left(b^{\prime}\right)-\varphi(b)\right)$ and $\varphi_{i}\left(b \otimes b^{\prime}\right)=\max \left(\varphi_{i}\left(b^{\prime}\right), \varphi_{i}(b)+\varphi_{i}\left(b^{\prime}\right)-\varepsilon\left(b^{\prime}\right)\right)$ are also used. The decomposition of such a tensor product is given by applying (7) repeatedly.

## 4. Crystal isomorphism and the symmetric group

Now let $\lambda_{1}$ and $\lambda_{2}$ be rectangular shape Young diagrams. Then $m_{\nu}=1$ for any $v$ in (7) hence we have a unique isomorphism of crystals $B_{\lambda_{1}} \otimes B_{\lambda_{2}} \simeq B_{\lambda_{2}} \otimes B_{\lambda_{1}}$ [SW, S$]$ which is given by tableaux products $[\mathrm{Fl}]$.

If $\lambda_{1}=\lambda_{2}$ the isomorphism is trivial, i.e. it is given by the identity map. Below we give explicit forms for non-trivial isomorphism between $B_{\ell}=B_{(\ell)}, B_{1}=B_{(1)}$ and $B_{\natural}=B_{(1,1)}$, which will be used in the description of the basic coloured system.

Let $\iota: B_{\ell} \otimes B_{1} \xrightarrow{\sim} B_{1} \otimes B_{\ell}$ be the map for the $s l_{n}$ crystal isomorphism. The decomposition by LR rule is given by $B_{\ell} \otimes B_{1} \simeq B_{(\ell+1)} \oplus B_{(\ell, 1)}$ and we have

where $p$ is determined by the condition $\alpha_{p}<\beta \leqslant \alpha_{p+1}$. The upper (resp. lower) one is associated with $B_{(\ell+1)}\left(\right.$ resp. $\left.B_{(\ell, 1)}\right)$.

Let $\iota^{\prime}: B_{\natural} \otimes B_{1} \xrightarrow{\sim} B_{1} \otimes B_{\square}$ be the map for the $s l_{n}$ crystal isomorphism. The decomposition by LR rule is given by $B_{\natural} \otimes B_{1} \simeq B_{(2,1)} \oplus B_{(1,1,1)}$ and we have

The first and the second ones are associated with $B_{(2,1)}$, and the third one with $B_{(1,1,1)}$.
Remark 4.1. The three cases here correspond to the loading-unloading processes of the decoding carrier which are denoted by $e, f, g$ in (2). Below we sometimes call an element of $B_{\square}$ a carrier for this reason.

Let $\iota^{\prime \prime}: B_{\ell} \otimes B_{\natural} \xrightarrow{\sim} B_{\natural} \otimes B_{\ell}$ be map for the $s l_{n}$ crystal isomorphism. The decomposition by LR rule is given by $B_{\ell} \otimes B_{\natural} \simeq B_{(\ell+1,1)} \oplus B_{(\ell, 1,1)}$ and we have

The first and the second ones are associated with $B_{(\ell, 1,1)}$, and the third and fourth ones with $B_{(\ell+1,1)}$. We shall omit explicit expressions for $\iota^{-1},\left(\iota^{\prime}\right)^{-1}$ but give one for $\left(\iota^{\prime \prime}\right)^{-1}$ in section 6 .

Now we consider a tensor product of $\mathcal{N}$ crystals $\mathcal{B}:=B_{\lambda_{1}} \otimes \cdots \otimes B_{\lambda_{\mathcal{N}}}$ where each $\lambda_{i}$ $(1 \leqslant i \leqslant \mathcal{N})$ is a rectangle. We let $\sigma$ denote the map for the isomorphism between tensor products of any two such crystals of rectangular shape Young tableaux.

Definition 4.2. For $\mathbf{p}=b_{1} \otimes \cdots \otimes b_{\mathcal{N}} \in \mathcal{B}$ we define $\sigma_{i}(1 \leqslant i \leqslant \mathcal{N}-1)$ by

$$
\sigma_{i}(\mathbf{p})=b_{1} \otimes \cdots \otimes \sigma\left(b_{i} \otimes b_{i+1}\right) \otimes \cdots \otimes b_{\mathcal{N}}
$$

Then we have
Proposition 4.3. Fix a positive integer $\ell$ and consider a tensor product $\mathcal{B}=B_{\lambda_{1}} \otimes \cdots \otimes B_{\lambda_{\mathcal{N}}}$ with each $\lambda_{i}(1 \leqslant i \leqslant \mathcal{N})$ being one of $(\ell),(1),(1,1)$. Then the $\sigma_{i}$ 's generate the symmetric group, i.e.

$$
\begin{aligned}
& \sigma_{i}^{2}=\mathrm{Id} \\
& \sigma_{i} \sigma_{j}=\sigma_{j} \sigma_{i} \text { for }|i-j| \geqslant 2 \\
& \sigma_{i} \sigma_{i+1} \sigma_{i}=\sigma_{i+1} \sigma_{i} \sigma_{i+1}
\end{aligned}
$$

Proof. Since $\sigma_{i}$ is the unique isomorphism the first identity holds. The second identity is obvious. We prove the third (Yang-Baxter) identity. Consider any such $\mathcal{B}=B_{\lambda_{1}} \otimes B_{\lambda_{2}} \otimes B_{\lambda_{3}}$ where $\lambda_{1}, \lambda_{2}, \lambda_{3}$ are all different. It is sufficient to check the identity on this $\mathcal{B}$. Moreover we can set $\lambda_{1}=(\ell), \lambda_{2}=(1)$ and $\lambda_{3}=(1,1)$ since the other cases are derived from it, by applying suitable compositions of $\sigma_{1}$ and $\sigma_{2}$. Let $b_{1} \otimes b_{2} \otimes b_{3}$ be any element of $\mathcal{B}$. There are an $s l_{n}$ highest weight element $u$ in $\mathcal{B}$ and a sequence $i_{1}, \ldots, i_{m}$ for some $m \in \mathbb{Z} \geqslant 0$ such that $b_{1} \otimes b_{2} \otimes b_{3}=\tilde{f}_{i_{m}} \ldots \tilde{f}_{i_{1}} u$. Therefore it is sufficient to prove

$$
\begin{equation*}
\sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} u=u \tag{8}
\end{equation*}
$$

for every $s l_{n}$ highest weight element in $\mathcal{B}$, since then we have $\sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1}\left(b_{1} \otimes b_{2} \otimes b_{3}\right)=$ $\sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} \tilde{f}_{i_{m}} \ldots \tilde{f}_{i_{1}} u=\tilde{f}_{i_{m}} \ldots \tilde{f}_{i_{1}} \sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} u=\tilde{f}_{i_{m}} \ldots \tilde{f}_{i_{1}} u=b_{1} \otimes b_{2} \otimes b_{3}$. By the LR rule [ $\mathrm{Fl}, \mathrm{Sa}$ ] we have
$B_{(\ell)} \otimes B_{(1)} \otimes B_{(1,1)}=B_{(\ell, 1,1,1)} \oplus B_{(\ell, 2,1)} \oplus B_{(\ell+1,2)} \oplus B_{(\ell+2,1)} \oplus\left(B_{(\ell+1,1,1)}\right)^{\oplus 2}$.
Thus for $\lambda=(\ell, 1,1,1),(\ell, 2,1),(\ell+1,2)$ or $(\ell+2,1)$ its associated $s l_{n}$ highest weight element is unique and hence (8) follows. For $\lambda=(\ell+1,1,1)$ there are two highest weight elements, for which we can verify (8) directly (see appendix B).

Remark 4.4. According to $[\mathrm{S}, \mathrm{SW}]$ the Yang-Baxter identity holds for any three $s l_{n}$ crystals of rectangular shape Young tableaux, hence the above proposition holds with each $\lambda_{i}$ being a generic rectangle. In the above we gave an elementary proof of this identity in the special case required for us.

## 5. Main theorem for the basic case

The highest weight elements of $B_{\ell}$ and $B_{\square}$ are given by

$$
u_{\ell}=\begin{array}{|l|l|l|l|}
\hline 1 & 1 & \cdots & 1 \\
\hline
\end{array} \in B_{\ell}, \quad u_{\natural}=\begin{array}{|c|}
\hline 1 \\
\hline 2 \\
\hline
\end{array} \quad \in B_{\natural} .
$$

Define the set of basic paths by

$$
\mathcal{P}=\left\{\mathbf{p}=p_{1} \otimes p_{2} \otimes \cdots \in B_{1}^{\otimes \infty} \mid p_{i}=1 \text { for } i \gg 1\right\}
$$

A basic path is regarded as an infinite array of boxes of capacity one with finite number of balls scattered among them, where 1 represents an empty box and $\alpha(\alpha \geqslant 2)$ a box containing a ball with index $\alpha$. We adopted a set of half-infinite paths as the space for the states of the automaton, but the formulation here is essentially not different from that in section 2.

Let $B_{\sharp}^{\prime}$ be a subset of $B_{\square}$ whose every element has a ' 1 '. We introduce operators $T_{\ell}(\ell \geqslant 1)$ and $T_{\natural}$ on $\mathcal{P}$ as follows. For any $\mathbf{p} \in \mathcal{P}$ we define $T_{\ell}(\mathbf{p}) \in \mathcal{P}$ and $T_{\natural}(\mathbf{p}) \in \mathcal{P}$ by using the maps for the $s l_{n}$ crystal isomorphism as

$$
\begin{align*}
& u_{\ell} \otimes \mathbf{p} \stackrel{\sim}{\mapsto} T_{\ell}(\mathbf{p}) \otimes u_{\ell}, \\
& u_{\text {曰 }} \otimes \mathbf{p} \stackrel{\sim}{\mapsto} T_{\natural}(\mathbf{p}) \otimes b(\mathbf{p}) . \tag{9}
\end{align*}
$$

Here $b(\mathbf{p}) \in B_{\natural}^{\prime}$ depends on the path $\mathbf{p}$. The application of $T_{\natural}$ on a path is equivalent to the procedure conducted by the decoding carrier in section 2 , and we can regard $b(\mathbf{p})$ as an outgoing carrier which takes off a letter $\geqslant 2$. (See example 2.5.) In what follows we occasionally use this terminology in some cases, and in particular we say that the decoding carrier takes off a '2' (resp. a letter $\geqslant 3$ ) when $b(\mathbf{p})=u_{\natural}\left(\right.$ resp. $\left.b(\mathbf{p}) \neq u_{\sharp}\right)$.

The operator $T_{\ell}$ gives a time evolution of the automaton that can be described by a carrier of capacity $\ell$ [TM, FOY]. For $\ell=\infty$ we have

Proposition 5.1 ([FOY]). The operator $T_{\infty}$ gives the same time evolution of the basic coloured system as in definition 2.1.

For $\mathbf{p}=p_{1} \otimes p_{2} \otimes \cdots \in \mathcal{P}$ let $F=F(\mathbf{p})=\max \left\{i \mid p_{i} \neq 1\right\}$ be the position of the rightmost non-empty box. We shall describe the map $\iota^{\prime}$ in section 4 in terms of the processes depicted in (1), (2). It is easy to see that

Lemma 5.2. When $T_{\mathrm{q}}$ is applied on $\mathbf{p}$,

1. The possible process that occurs at the position $F+1$ is a or $b$ in (1).
2. The possible process that occurs at every position $\geqslant F+2$ is a in (1).

Then we have
Lemma 5.3. Apply $T_{\natural}$ on $\mathbf{p}$. If the decoding carrier takes off a ' 2 ' then $F\left(T_{\natural}(\mathbf{p})\right)=F(\mathbf{p})$.
Proof. From the above lemma the only possible process that occurs at $F+1$ is $a$ in (1) in this case.

Lemma 5.4. Suppose all the decoding carriers take off ' 2 's when $\left(T_{\square}\right)^{k}$ is applied on $\mathbf{p}$. Then there is no letter $\geqslant 3$ at the positions $\geqslant F-k+1$ in $\mathbf{p}$.

Proof. Induction on $k$. Suppose there is a letter $\geqslant 3$, say $\gamma$, at the position $F$. By applying $T_{\square}$ on $\mathbf{p}$, the possible process that occurs at $F$ is one of $c-g$ in (1), (2). Hence the $\gamma$ is loaded into the carrier in any case. Then by lemma 5.2 the decoding carrier should take off a letter $\geqslant 3$. (See the following example.)

Suppose there is a letter $\geqslant 3$, say $\gamma$, at the position $F-k+1$ and the decoding carrier takes off a ' 2 ' when $T_{\natural}$ is applied on $\mathbf{p}$. If so, this $\gamma$ should go rightwards by at least one box position at this time. Therefore there is at least one letter $\geqslant 3$ at some position $\geqslant F-k+2$ in $T_{\mathrm{t}}(\mathbf{p})$. The proof follows by induction and lemma 5.3.

Example 5.5. From (1) and (2) the possible loading-unloading processes which can occur at $F$ and $F+1$ are as follows.


Therefore if $\gamma \geqslant 3$ then the decoding carrier takes off a letter $\geqslant 3$.
By taking $k=F$ we have
Corollary 5.6. Suppose all the decoding carriers take off ' 2 's when $\left(T_{\square}\right)^{F}$ is applied on $\mathbf{p}$. Then there is no letter $\geqslant 3$ in $\mathbf{p}$.

From this corollary we can deduce that
Theorem 5.7. For any $\mathbf{p} \in \mathcal{P}$ all the letters $\geqslant 3$ in $\mathbf{p}$ can be removed by applying $T_{\natural}$ sufficiently many times.

Thus for any $\mathbf{p} \in \mathcal{P}$ we can find a non-negative integer $N$ such that both $\tilde{\mathbf{p}}:=\left(T_{\natural}\right)^{N}(\mathbf{p})$ and $\widetilde{T_{\ell}(\mathbf{p})}:=\left(T_{\square}\right)^{N}\left(T_{\ell}(\mathbf{p})\right)$ are composed of only 1 and 2 . Choose any integer $N$ such that these conditions are satisfied. By using (9) we have

$$
\begin{align*}
& u_{\natural}^{\otimes N} \otimes \mathbf{p} \stackrel{\sim}{\mapsto} \tilde{\mathbf{p}} \otimes\left(b_{1} \otimes \cdots \otimes b_{N}\right), \\
& u_{\natural}^{\otimes N} \otimes T_{\ell}(\mathbf{p}) \stackrel{\sim}{\mapsto} \widetilde{T_{\ell}(\mathbf{p})} \otimes\left(b_{1}^{\prime} \otimes \cdots \otimes b_{N}^{\prime}\right), \tag{10}
\end{align*}
$$

where $b_{i}, b_{i}^{\prime} \in B_{\natural}^{\prime}(1 \leqslant i \leqslant N)$ are defined as $b_{i}=b\left(\left(T_{\natural}\right)^{N-i}(\mathbf{p})\right)$ and $b_{i}^{\prime}=b\left(\left(T_{\text {घ }}\right)^{N-i}\left(T_{\ell}(\mathbf{p})\right)\right)$. It is easy to see that if $\mathbf{p} \in \mathcal{P}$ has no letter $\geqslant 3$ then $T_{\natural}(\mathbf{p})=\mathbf{p}$. Thus the $\tilde{\mathbf{p}}$ does not depend on the choice of $N$.

Now we present the main theorem in this paper.
Theorem 5.8. With the notation as above the following relations hold: $T_{\ell}(\tilde{\mathbf{p}})=\widetilde{T_{\ell}(\mathbf{p})}$ and $b_{i}=b_{i}^{\prime}$ for $1 \leqslant i \leqslant N$.

Proof. By the isomorphism of $s l_{n}$ crystal we have

$$
\begin{aligned}
u_{\ell} \otimes u_{\natural}^{\otimes N} \otimes \mathbf{p} & \stackrel{\sim}{\mapsto} u_{\natural}^{\otimes N} \otimes u_{\ell} \otimes \mathbf{p} \\
& \stackrel{\sim}{\mapsto} u_{\sharp}^{\otimes N} \otimes T_{\ell}(\mathbf{p}) \otimes u_{\ell} \\
& \stackrel{\sim}{\mapsto} \widetilde{T_{\ell}(\mathbf{p})} \otimes\left(b_{1}^{\prime} \otimes \cdots \otimes b_{N}^{\prime}\right) \otimes u_{\ell},
\end{aligned}
$$

and

$$
\begin{aligned}
u_{\ell} \otimes u_{\natural}^{\otimes N} \otimes \mathbf{p} & \stackrel{\sim}{\mapsto} u_{\ell} \otimes \tilde{\mathbf{p}} \otimes\left(b_{1} \otimes \cdots \otimes b_{N}\right) \\
& \stackrel{\sim}{\mapsto} T_{\ell}(\tilde{\mathbf{p}}) \otimes u_{\ell} \otimes\left(b_{1} \otimes \cdots \otimes b_{N}\right) \\
& \stackrel{\sim}{\mapsto} T_{\ell}(\tilde{\mathbf{p}}) \otimes\left(b_{1} \otimes \cdots \otimes b_{N}\right) \otimes u_{\ell} .
\end{aligned}
$$

For any positive integer $L$ we consider a tensor product of crystals $\mathcal{B}=B_{\ell} \otimes B_{\square}^{\otimes N} \otimes B_{1}^{\otimes L} \simeq$ $B_{1}^{\otimes L} \otimes B_{\natural}^{\otimes N} \otimes B_{\ell}$. For this $\mathcal{B}$ define $\sigma_{i}$ as in definition 4.2 with $\mathcal{N}=L+N+1$. To prove the theorem it is sufficient to show the following $x$ and $y$ coincide:

$$
\begin{aligned}
& x:=\left(\sigma_{L} \ldots \sigma_{2} \sigma_{1}\right)\left(\sigma_{L+1} \ldots \sigma_{3} \sigma_{2}\right) \ldots\left(\sigma_{N+L-1} \ldots \sigma_{N+1} \sigma_{N}\right)\left(\sigma_{N+L} \ldots \sigma_{2} \sigma_{1}\right) \\
& y:=\left(\sigma_{N+L} \ldots \sigma_{2} \sigma_{1}\right)\left(\sigma_{L+1} \ldots \sigma_{3} \sigma_{2}\right)\left(\sigma_{L+2} \ldots \sigma_{4} \sigma_{3}\right) \ldots\left(\sigma_{N+L} \ldots \sigma_{N+2} \sigma_{N+1}\right)
\end{aligned}
$$

They serve as two possible ways to send an element of $B_{\ell} \otimes B_{\natural}^{\otimes N} \otimes B_{1}^{\otimes L}$ into $B_{1}^{\otimes L} \otimes B_{\natural}^{\otimes N} \otimes B_{\ell}$. The $x$ is for $B_{\ell} \otimes B_{\natural}^{\otimes N} \otimes B_{1}^{\otimes L} \xrightarrow{\sim} B_{\natural}^{\otimes N} \otimes B_{\ell} \otimes B_{1}^{\otimes L} \xrightarrow{\sim} B_{\square}^{\otimes N} \otimes B_{1}^{\otimes L} \otimes B_{\ell} \xrightarrow{\sim} B_{1}^{\otimes L} \otimes B_{\natural}^{\otimes N} \otimes B_{\ell}$ and the $y$ is for $B_{\ell} \otimes B_{\square}^{\otimes N} \otimes B_{1}^{\otimes L} \xrightarrow{\sim} B_{\ell} \otimes B_{1}^{\otimes L} \otimes B_{\natural}^{\otimes N} \xrightarrow{\sim} B_{1}^{\otimes L} \otimes B_{\ell} \otimes B_{\natural}^{\otimes N} \xrightarrow{\sim} B_{1}^{\otimes L} \otimes B_{\natural}^{\otimes N} \otimes B_{\ell}$. By proposition 4.3 one can verify $x=y$ directly. (See the following example.)

Example 5.9. Let $N=L=3$. For notational simplicity we denote $\sigma_{i}$ by $i$. By repeated use of $\sigma_{i} \sigma_{i+1} \sigma_{i}=\sigma_{i+1} \sigma_{i} \sigma_{i+1}$ we have $1234321=4321234$ for instance. Then

$$
\begin{aligned}
(321)(432)(543)(654321) & =(32)(43)(54)(654321)(234) \\
& =(\mathbf{3})(\mathbf{4})(\mathbf{5})(\mathbf{6 5 4 3 2 1})(345)(234) \\
& =(654321)(456)(345)(234) \\
& =(654321)(432)(543)(654) .
\end{aligned}
$$

Here we used $\sigma_{i} \sigma_{j}=\sigma_{j} \sigma_{i}$ for $|i-j| \geqslant 2$.
Let $y_{i}(\neq 1)$ be the letter that appears in the $b_{i}(1 \leqslant i \leqslant N)$ in (10). Denote the word $y_{1} \ldots y_{N}$ by $\mathbf{y}$ and write the relation (10) symbolically as $\mathbf{p}=\tilde{\mathbf{p}} \oplus \mathbf{y}$. Then by theorem 5.8 we have

$$
\begin{equation*}
T_{\ell}(\mathbf{p})=T_{\ell}(\tilde{\mathbf{p}}) \oplus \mathbf{y} \tag{11}
\end{equation*}
$$

for any $\ell$. Thus by letting $\ell$ be infinity we obtain (4) by proposition 5.1.

## 6. Generalization to the inhomogeneous case

For the inhomogeneous case we need to prove the Yang-Baxter identity on $B_{\ell_{1}} \otimes B_{\ell_{2}} \otimes B_{\ell_{3}}$ and on $B_{\ell_{1}} \otimes B_{\ell_{2}} \otimes B_{\natural}$. Although there is a more general assertion mentioned in remark 4.4, we shall still give some explanations on this identity in these special cases in order to make our arguments self-contained.

As a set the $s l_{n}$ crystal $B_{\ell}$ in (5) is also given by

$$
\begin{equation*}
B_{\ell}=\left\{\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{Z}_{\geqslant 0}^{n} \mid \sum_{i=1}^{n} x_{i}=\ell\right\} \tag{12}
\end{equation*}
$$

Here $x_{i}$ shows the number of the letter $i$ that appears in the tableau representation (5). Let $B$ and $B^{\prime}$ be any two of these $s l_{n}$ crystals. Given $x=\left(x_{1}, \ldots, x_{n}\right) \in B, y=\left(y_{1}, \ldots, y_{n}\right) \in B^{\prime}$ we define a piecewise linear map $R:(x, y) \mapsto\left(x^{\prime}, y^{\prime}\right)$ as $x^{\prime}=\left(x_{1}^{\prime}, \ldots, x_{n}^{\prime}\right), y^{\prime}=\left(y_{1}^{\prime}, \ldots, y_{n}^{\prime}\right)$ and $x_{i}^{\prime}=y_{i}+P_{i+1}-P_{i}, y_{i}^{\prime}=x_{i}+P_{i}-P_{i+1}$, where $P_{i}=\max _{1 \leqslant j \leqslant n}\left(\sum_{k=1}^{j-1}\left(y_{k+i-1}-x_{k+i-1}\right)+\right.$ $\left.y_{j+i-1}\right)$. Here the indices are interpreted in modulo $n$. This piecewise linear map gives the isomorphism of $s l_{n}$ crystals $B \otimes B^{\prime} \xrightarrow{\sim} B^{\prime} \otimes B$ [HHIKTT]. Consider such $\mathcal{B}=B_{\ell_{1}} \otimes B_{\ell_{2}} \otimes B_{\ell_{3}}$. For this $\mathcal{B}$ define $\sigma_{i}$ as in definition 4 with $\mathcal{N}=3$ and $\sigma=R$. Then the Yang-Baxter identity $\sigma_{1} \sigma_{2} \sigma_{1}=\sigma_{2} \sigma_{1} \sigma_{2}$ holds on this $\mathcal{B}$. A simple proof of this identity on this $\mathcal{B}$ is given by using a birational map analogue of the above $R$ [KNY, Y].

Next we consider such $\mathcal{B}=B_{\ell_{1}} \otimes B_{\ell_{2}} \otimes B_{\natural}$. Define $\sigma_{i}$ as above but now $\sigma$ is one of $R, \iota^{\prime \prime},\left(\iota^{\prime \prime}\right)^{-1}$.

Lemma 6.1. The identity $\sigma_{1} \sigma_{2} \sigma_{1}=\sigma_{2} \sigma_{1} \sigma_{2}$ holds on $\mathcal{B}=B_{\ell_{1}} \otimes B_{\ell_{2}} \otimes B_{\natural}$.
Proof. For simplicity we assume $\ell_{1}>\ell_{2}$. It is sufficient to prove $\sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} u=u$ for every $s l_{n}$ highest weight element $u$ in $\mathcal{B}$. By the LR rule [Fl, Sa] we have

$$
\begin{aligned}
B_{\left(\ell_{1}\right)} \otimes B_{\left(\ell_{2}\right)} \otimes B_{(1,1)}= & \left(\bigoplus_{x=1}^{\ell_{2}} B_{\left(\ell_{1}+\ell_{2}-x, x, 1,1\right)}\right) \oplus\left(\bigoplus_{x=0}^{\ell_{2}} B_{\left(\ell_{1}+\ell_{2}-x+1, x+1\right)}\right) \\
& \oplus\left(\bigoplus_{x=0}^{\ell_{2}-1}\left(B_{\left(\ell_{1}+\ell_{2}-x, x+1,1\right)}\right)^{\oplus 2}\right) \oplus B_{\left(\ell_{1}, \ell_{2}+1,1\right)} .
\end{aligned}
$$

So it is enough to prove the above identity only for the highest weight elements associated with $B_{\left(\ell_{1}+\ell_{2}-x, x+1,1\right)}$ for $0 \leqslant x \leqslant \ell_{2}-1$ because the other $B_{\lambda} \mathrm{s}$ are multiplicity free. It can be verified directly. (See appendix C.)

For a given sequence of positive integers $\ell=\left(\ell_{1}, \ell_{2}, \ldots\right)$ define the associated set of inhomogeneous paths by

$$
\mathcal{P}_{\ell}=\left\{\mathbf{p}=p_{1} \otimes p_{2} \otimes \cdots \in B_{\ell_{1}} \otimes B_{\ell_{2}} \otimes \cdots \mid p_{i}=u_{\ell_{i}} \text { for } i \gg 1\right\} .
$$

An inhomogeneous path is regarded as an infinite array of boxes of various capacities with finite number of balls scattered among them. At the $i$ th position there is a box $p_{i}$ of capacity $\ell_{i}$. If it has the expression $p_{i}=\left(x_{1}, \ldots, x_{n}\right)$ then we interpret it as a box containing $x_{k}$ balls with index $k$ for $2 \leqslant k \leqslant n$. We define the operators $T_{\ell}(\ell \geqslant 1)$ and $T_{\text {曰 }}$ which act on $\mathcal{P}_{\ell}$ by using the same formulae in (9). As in the basic case the operator $T_{\ell}$ gives the time evolution described by a carrier of capacity $\ell$ [HHIKTT]. For $\ell=\infty$ the operator $T_{\infty}$ also has a description that is a generalization of definition 2.1 to the inhomogeneous case [TTM].

We gave an explicit formula for the map $\iota^{\prime \prime}: B_{\ell} \otimes B_{\natural} \xrightarrow{\sim} B_{\natural} \otimes B_{\ell}$ in section 4. The inverse of $\iota^{\prime \prime}$ is given as follows


This formula shows the loading-unloading processes of the decoding carrier at a box with capacity $\geqslant 2$, corresponding to those pictures in (1), (2) for the case of capacity one. As in the basic case let $F$ be the position of the rightmost non-empty box i.e. $F=F(\mathbf{p})=\max \left\{i \mid p_{i} \neq u_{\ell_{i}}\right\}$ for $\mathbf{p}=p_{1} \otimes p_{2} \otimes \cdots \in \mathcal{P}_{\ell}$. It is easy to see that

Lemma 6.2. When $T_{\natural}$ is applied on $\mathbf{p}$,

1. The possible process that occurs at the position $F+1$ is (I).
2. The possible process that occurs at every position $\geqslant F+2$ is $(I)$ with $\alpha=1$.

Then we have
Lemma 6.3. Apply $T_{\natural}$ on $\mathbf{p}$. If the decoding carrier takes off a ' 2 ' then $F\left(T_{\natural}(\mathbf{p})\right)=F(\mathbf{p})$.
Proof. From the above lemma the only possible process that occurs at $F+1$ is (I) with $\alpha=1$ in this case.

Lemma 6.4. Suppose all the decoding carriers take off ' 2 's when $\left(T_{\sharp}\right)^{\ell_{F}+\ell_{F-1}+\cdots+\ell_{F-k+1}}$ is applied on $\mathbf{p}$. Then there is no letter $\geqslant 3$ at the positions $\geqslant F-k+1$ in $\mathbf{p}$.

Proof. Induction on $k$. Suppose all the decoding carriers take off ' 2 's when $\left(T_{\natural}\right)^{\ell_{F}}$ is applied on $\mathbf{p}$. If so, no carrier that leaves from the position $F$ has a letter $\geqslant 3$, by lemma 6.2. This implies that under the application of $\left(T_{t}\right)^{\ell_{F}}$ the possible processes that occur at the position $F$ are (I), (II), (III), and not (IV) or (V). Now suppose there is a letter $\geqslant 3$, say $\gamma$, in the tableau at the position $F$. In (I), (II), (III) any letter in the tableau is either loaded into the carrier or shifted to the left inside the tableau. In particular the leftmost letter $\gamma_{1}$ is always loaded into the carrier. Therefore the $\gamma$ should be loaded into some carrier when $\left(T_{\square}\right)^{\ell_{F}}$ is applied on $\mathbf{p}$, because the width of the tableau is $\ell_{F}$. This contradicts the above observation. Thus there is no letter $\geqslant 3$ at the position $F$.

Suppose there is a letter $\geqslant 3$, say $\gamma$, at the position $F-k+1$ and all the decoding carriers take off ' 2 's when $\left(T_{\mathrm{t}}\right)^{\ell_{F-k+1}}$ is applied on $\mathbf{p}$. Suppose also that the processes that occur at the position $F-k+1$ are (I), (II), (III), and not (IV) or (V) at this time. If so, the $\gamma$ should be loaded into some carrier when $\left(T_{\square}\right)^{\ell_{F-k+1}}$ is applied on $\mathbf{p}$, by the same reason in the previous paragraph. If not, there should be at least one letter $\geqslant 3$ delivered to some position $\geqslant F-k+2$ by a carrier which has left from the position $F-k+1$ after the process (IV) or (V). In any case there should be at least one letter $\geqslant 3$ at some position $\geqslant F-k+2$ in $\left(T_{\sharp}\right)^{\ell_{F-k+1}}(\mathbf{p})$. The proof follows by induction and lemma 6.3.

By taking $k=F$ we have
Corollary 6.5. Suppose all the decoding carriers take off ' 2 's when $\left(T_{\natural}\right)^{\ell_{F}+\ell_{F-1}+\cdots+\ell_{1}}$ is applied on $\mathbf{p}$. Then there is no letter $\geqslant 3$ in $\mathbf{p}$.

Therefore
Lemma 6.6. For any $\mathbf{p} \in \mathcal{P}_{\ell}$ all the letters $\geqslant 3$ in $\mathbf{p}$ can be removed by applying $T_{\square}$ sufficiently many times.

By this lemma and the Yang-Baxter identity we can deduce that the separation of colour degree of freedom (theorem 5.8) also holds in the inhomogeneous case.

## 7. Concluding remarks

There is a clear notion of Liouville integrability for continuous dynamical Hamiltonian systems with finite degrees of freedom. Since the box-ball systems are (ultra-)discrete dynamical systems with infinite degrees of freedom, the notion of integrability (or completeness of conserved quantities) is not so clear. Nevertheless the author thinks it reasonable to say that all the conserved quantities of the basic coloured system have been obtained by the method in this paper, by combining with that in [Tkg].

The reason is as follows. What we have done is not a mere construction of conserved quantities, but an explicit dynamical variable change which makes part of the variables ( $\mathbf{y}$ in (3)) a conserved quantity. Then by applying the method in [Tkg] to the underlying monochrome system (to which $\tilde{\mathbf{p}}$ in (3) belongs), the coloured system can be transformed into a set of non-interacting linear dynamical systems in which the time evolution is trivial. We note that the variable change is invertible because the inverse of $T_{\natural}$ can be defined if $b(\mathbf{p})$ in (9) is given. It allows us to write out any solution of the initial value problems of the original system by solutions of those linear dynamical systems and a word $\mathbf{y}$. In this sense a complete set of conserved quantities of the system consists of the conserved quantities of the linear dynamical systems and the word.

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## Appendix A. An example of $\boldsymbol{T}(\mathbf{p})=T(\tilde{\mathbf{p}}) \oplus \mathbf{y}$

For $t=1$ of example 2.4 we obtain the following separation:


For $t=2$ we obtain the following:


For $t=3$ we obtain the following:

| $\mathrm{s}=0$ | 435. 54222 . . . . . 2 |
| :---: | :---: |
| $\mathrm{s}=1$ | 5423.55422....... 4 |
| $\mathrm{s}=2$ | 5242.55322..... 5 |
| $\mathrm{s}=3$ | . 2522.54322 ..... 2 |
| $\mathrm{s}=4$ | . 2.22255432 . . . 3 |
| $\mathrm{s}=5$ | . $2.222 .55422 . . .4$ |
| $\mathrm{s}=6$ | . 2 . 222 . 55222 . . 5 |
| $s=7$ | . 2 . 222 . . 52222 . 5 |
| $s=8$ | 2. 222 . . . 22222. |

Note that we always obtain the same sequence of removed letters $\mathbf{y}=55432542$ (in reverse order), and that the automaton states in the last rows coincide with the states in example 2.3.

## Appendix B. Proof of proposition 4.3 (Continued)

For each highest element $u$ of the shape $\lambda=(\ell+1,1,1)$ the relation $\sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} u=u$ can be verified as (we have set $\ell=3$ here)

$$
\begin{aligned}
& \begin{array}{|l|l|l|}
\hline 1 & 1 & 1 \\
\hline
\end{array} \otimes \begin{array}{|l|}
\hline 2 \\
\hline 3 \\
\hline
\end{array}
\end{aligned}
$$

$$
\begin{aligned}
& \begin{array}{|l|l|l|}
\hline 1 & 1 & 1 \\
\hline
\end{array} \otimes \begin{array}{|l|}
\hline 1 \\
\hline 3 \\
\hline
\end{array}
\end{aligned}
$$

$$
\begin{aligned}
& \stackrel{\sim}{\mapsto} \begin{array}{|c|}
\hline 1 \\
\hline 2 \\
\hline
\end{array} \otimes \begin{array}{|l|l|l|}
\hline 1 & 1 & 3 \\
\hline
\end{array} \stackrel{\sim}{\mapsto} \begin{array}{|l|l|l|}
\hline 1 & 1 & 1 \\
\hline
\end{array} \otimes \begin{array}{|c|}
\hline 2 \\
\hline 3 \\
\hline 1 \\
\mapsto
\end{array} \begin{array}{|l|l|l|}
\hline 1 & 1 & 1 \\
\hline
\end{array} \otimes \begin{array}{|l|l|}
\hline 1 \\
\hline 3 \\
\hline
\end{array} .
\end{aligned}
$$

## Appendix C. Proof of lemma 6.1 (continued)

Recall the description of the elements of $B_{\ell}$ in (12). Denote $\left(x_{1}, x_{2}, x_{3}, 0, \ldots, 0\right)$ by [ $x_{1}, x_{2}, x_{3}$ ]. For each highest element $u$ of the shape $\lambda=\left(\ell_{1}+\ell_{2}-x, x+1,1\right)$ the relation $\sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} \sigma_{2} \sigma_{1} u=u$ can be verified as

$\left[\ell_{1}, 0,0\right] \otimes\left[\ell_{2}-x, x, 0\right] \otimes$| 2 |
| :--- |
| 3 |

$$
\begin{aligned}
& \stackrel{\sim}{\mapsto}\left[\ell_{2}, 0,0\right] \otimes\left[\ell_{1}-x, x, 0\right] \otimes \frac{2}{3} \quad \stackrel{\sim}{\mapsto}\left[\ell_{2}, 0,0\right] \otimes \begin{array}{|c|}
\hline 1 \\
\hline 2 \\
\hline
\end{array} \otimes\left[\ell_{1}-x-1, x, 1\right] \\
& \stackrel{\sim}{\mapsto}\left[\begin{array} { | c } 
{ 1 } \\
{ 2 }
\end{array} \otimes [ \ell _ { 2 } , 0 , 0 ] \otimes [ \ell _ { 1 } - x - 1 , x , 1 ] \stackrel { \sim } { \mapsto } \left[\begin{array}{|c}
1 \\
\hline 2
\end{array} \otimes\left[\ell_{1}, 0,0\right] \otimes\left[\ell_{2}-x-1, x, 1\right]\right.\right. \\
& \stackrel{\sim}{\mapsto}\left[\ell_{1}, 0,0\right] \otimes \underset{\frac{1}{2}}{\square} \otimes\left[\ell_{2}-x-1, x, 1\right] \stackrel{\sim}{\mapsto}\left[\ell_{1}, 0,0\right] \otimes\left[\ell_{2}-x, x, 0\right] \otimes \frac{2}{2}, \\
& {\left[\ell_{1}, 0,0\right] \otimes\left[\ell_{2}-x-1, x+1,0\right] \otimes \begin{array}{|l|}
\hline 1 \\
\hline 3 \\
\hline
\end{array}} \\
& \stackrel{\sim}{\mapsto}\left[\ell_{2}, 0,0\right] \otimes\left[\ell_{1}-x-1, x+1,0\right] \otimes \underset{\mid}{4} \stackrel{\sim}{\mapsto}\left[\ell_{2}, 0,0\right] \otimes \frac{2}{3} \otimes\left[\ell_{1}-x, x, 0\right] \\
& \stackrel{\sim}{\mapsto}\left[\begin{array} { l } 
{ 1 } \\
{ 2 }
\end{array} \otimes [ \ell _ { 2 } - 1 , 0 , 1 ] \otimes [ \ell _ { 1 } - x , x , 0 ] \stackrel { \sim } { \mapsto } \left[\begin{array}{l}
1 \\
\hline 2
\end{array} \otimes\left[\ell_{1}-1,0,1\right] \otimes\left[\ell_{2}-x, x, 0\right]\right.\right. \\
& \stackrel{\sim}{\mapsto}\left[\ell_{1}, 0,0\right] \otimes \begin{array}{|c}
\frac{2}{3}
\end{array} \otimes\left[\ell_{2}-x, x, 0\right] \stackrel{\sim}{\mapsto}\left[\ell_{1}, 0,0\right] \otimes\left[\ell_{2}-x-1, x+1,0\right] \otimes \frac{1}{3} .
\end{aligned}
$$

## References

[F] Fukuda K 2004 Box-ball systems and Robinson-Schensted-Knuth correspondence J. Algebr. Comb. 19 67-89
[FOY] Fukuda K, Okado M and Yamada Y 2000 Energy functions in box ball systems Int. J. Mod. Phys. A 15 1379-92
[Fl] Fulton W 1997 Young tableaux: with applications to representation theory and geometry London Math. Soc. Student Texts 35 (Cambridge: Cambridge University Press)
[HHIKTT] Hatayama G, Hikami K, Inoue R, Kuniba A, Takagi T and Tokihiro T 2001 The $A_{M}^{(1)}$ Automata related to crystals of symmetric tensors J. Math. Phys. 42 274-308
[HKOTY] Hatayama G, Kuniba A, Okado M, Takagi T and Yamada Y 2002 Scattering rules in soliton cellular automata associated with crystal bases Contemp. Math. 297 151-82
[KMN] Kang S-J, Kashiwara M, Misra K C, Miwa T, Nakashima T and Nakayashiki A 1992 Perfect crystals of quantum affine Lie algebras Duke Math. J. 68 499-607
[KNY] Kajiwara K, Noumi M and Yamada Y 2002 Discrete dynamical systems with $W\left(A_{m-1}^{(1)} \times A_{n-1}^{(1)}\right)$ symmetry Lett. Math. Phys. 60 211-9
[K1] Kashiwara M 1990 Crystallizing the $q$-analogue of universal enveloping algebras Commun. Math. Phys. 133 249-60
[K2] Kashiwara M 1994 Crystal bases of modified quantized enveloping algebra Duke Math. J. 73 383-413
[KN] Kashiwara M and Nakashima T 1994 Crystal graph for representations of the q-analogue of classical Lie algebras J. Algebra 165 295-345
[N] Nakashima T 1993 Crystal base and a generalization of Littlewood-Richardson rule for the classical Lie algebras Commun. Math. Phys. 154 215-43
[Sa] Sagan B E 2001 The symmetric group: representations, combinatorial algorithms, and symmetric functions Graduate Texts in Mathematics 2nd edn 203 (Berlin: Springer)
[SW] Schilling A and Warnaar S O 1999 Inhomogeneous lattice paths, generalized Kostka polynomials and $A_{n}-1$ supernomials Commun. Math. Phys. 202 359-401
[S] Shimozono M 2002 Affine type A crystal structure on tensor product of rectangles, demazure characters, and nilpotent varieties J. Algebr. Comb. 15 151-87
[Tkg] Takagi T 2005 Inverse scattering method for a soliton cellular automaton Nucl. Phys. B 707 577-601 (Preprint math-ph/0406038) at press
[T] Takahashi D 1993 On some soliton systems defined by using boxes and balls Proc. Int. Symp. on Nonlinear Theory and Its Applications (NOLTA '93) pp 555-8
[TM] Takahashi D and Matsukidaira J 1997 Box and ball system with a carrier and ultra-discrete modified KdV equation J. Phys. A: Math. Gen. 30 L733-9
[TS] Takahashi D and Satsuma J 1990 A soliton cellular automaton J. Phys. Soc. Japan. 59 3514-9
[TNS] Tokihiro T, Nagai A and Satsuma J 1999 Proof of the solitonical nature of box and ball systems by means of inverse ultra-discretization Inverse Problems 15 1639-62
[TTM] Tokihiro T, Takahashi D and Matsukidaira J 2000 Box and ball system as a realization of ultradiscrete nonautonomous KP equation J. Phys. A: Math. Gen. 33 607-19
[TTS] Torii M, Takahashi D and Satsuma J 1996 Combinatorial representation of invariants of a soliton cellular automaton Physica D 92 209-20
[Y] Yamada Y 2001 A birational representation of Weyl group, combinatorial $R$-matrix and discrete Toda equation Physics and Combinatorics 2000 ed A N Kirillov and N Liskova Proc. Nagoya 2000 Int. Workshop (Singapore: World Scientific) pp 305-19

